

Introduction to quantum thermodynamics

An open systems perspective

Dr Graeme Pleasance

*Quantum@SUN group
Department of Physics
Stellenbosch University*

NITheCS Mini-School
7 Feb, 2024



Lecture 1 Feb 7: **Background** – foundational quantum mechanics, open quantum systems, equilibrium descriptions.

Lecture 2 Feb 14: **Markovian master equations** – microscopic derivation, 1st and 2nd laws of thermodynamics.

Lecture 3 Feb 21: **Quantum thermal machines** – discrete heat engines and refrigerators, quantum Otto cycle.

Lecture 4 Feb 28: **Strong system-reservoir coupling** – mean force Gibbs state, nonequilibrium descriptions.

Background

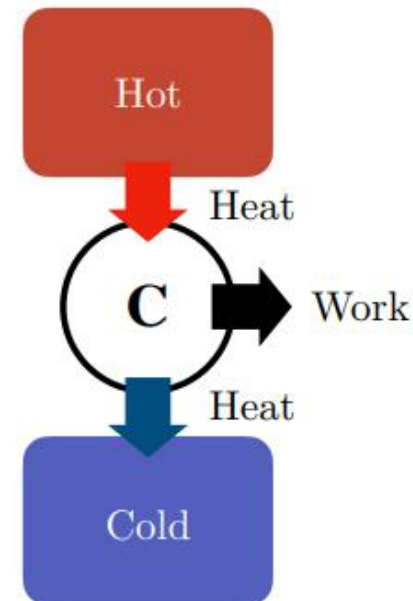
Macroscopic thermodynamics: phenomenological theory describing the *average* behavior of heat, work and entropy.

Quantum thermodynamics: extends 'classical' concepts of heat and work into the microscopic domain.

Background

Macroscopic thermodynamics: phenomenological theory describing the *average* behavior of heat, work and entropy.

Quantum thermodynamics: extends 'classical' concepts of heat and work into the microscopic domain.

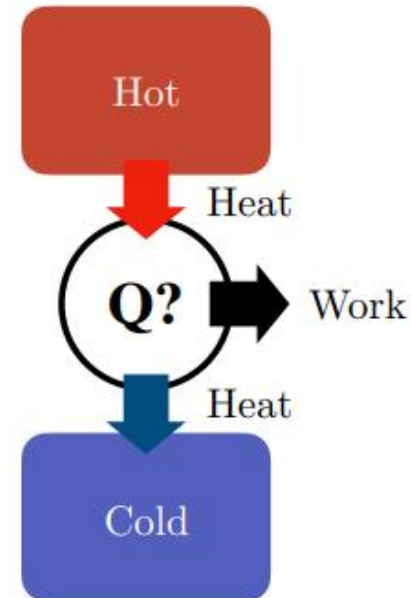


bluebell-railway.com/the_flying_scotsman_in_steam_xx/

Background

Macroscopic thermodynamics: phenomenological theory describing the *average* behavior of heat, work and entropy.

Quantum thermodynamics: extends 'classical' concepts of heat and work into the microscopic domain.



bluebell-railway.com/the_flying_scotsman_in_steam_xx/

Background

Macroscopic thermodynamics: phenomenological theory describing the *average* behavior of heat, work and entropy.

Quantum thermodynamics: extends 'classical' concepts of heat and work into the microscopic domain.

- Thermal machines.
- Thermalization.
- Stochastic thermodynamics (fluctuations).

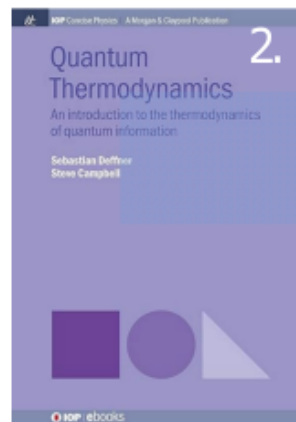
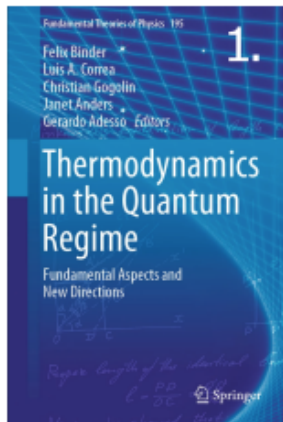
Background

Macroscopic thermodynamics: phenomenological theory describing the *average* behavior of heat, work and entropy.

Quantum thermodynamics: extends 'classical' concepts of heat and work into the microscopic domain.

- Thermal machines.
- Thermalization.
- Stochastic thermodynamics (fluctuations).

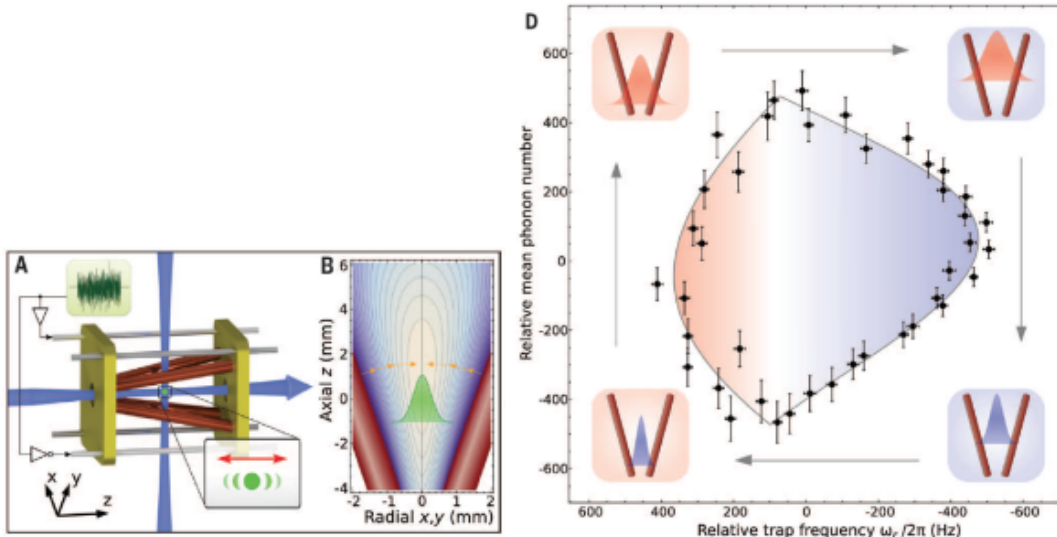
Useful literature:



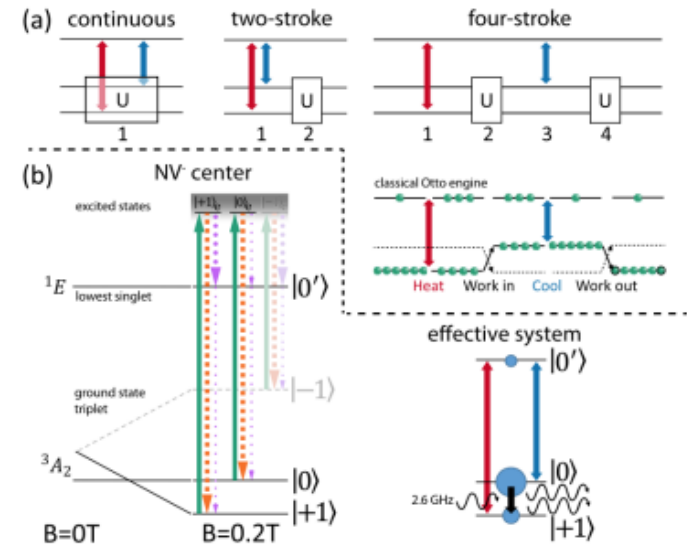
3. S. Vinjanampathy. & J. Anders, *Quantum thermodynamics*, Contemp. Phys, **57**, 545, (2016)
4. P. Potts, *Introduction to Quantum Thermodynamics*, arXiv:1906.07439 (2019)

Background

Technological applications:



(Roßnagel et al., Science **352**, 325-329, 2016)



(Klatzow et al., PRL, **122**, 110601, 2019)

Powering an Engine with Quantum Coherence

Janet Anders

Department of Physics and Astronomy, University of Exeter, Exeter, United Kingdom

March 20, 2019 • Physics 12, 32

Contents – lecture 1

1 Motivation

2 Isolated quantum systems

- Time evolution
- Spectral theorem
- Density matrix

3 Open quantum systems

- Reduced density matrix
- Quantum maps (Kraus representation)

4 Equilibrium descriptions

5 Summary

Preliminaries

Quantum states

State ψ of a **Isolated quantum system** represented by $|\psi\rangle \in \mathcal{H}$:

$$\mathcal{H} = \{|\psi\rangle = \begin{pmatrix} c_0 \\ \vdots \\ c_{d-1} \end{pmatrix} \mid c_i \in \mathbb{C}\}$$

Inner (scalar) product of $\phi, \psi \in \mathcal{H}$:

$$\langle \phi | \psi \rangle = \sum_k b_k^* c_k, \quad \langle \phi | = (b_0^*, \dots, b_{d-1}^*),$$

with $\langle \phi | \in \mathcal{H}^*$.

Preliminaries

Quantum states

State ψ of a **Isolated quantum system** represented by $|\psi\rangle \in \mathcal{H}$:

$$\mathcal{H} = \{|\psi\rangle = \begin{pmatrix} c_0 \\ \vdots \\ c_{d-1} \end{pmatrix} \mid c_i \in \mathbb{C}\}$$

Inner (scalar) product of $\phi, \psi \in \mathcal{H}$:

$$\langle \phi | \psi \rangle = \sum_k b_k^* c_k, \quad \langle \phi | = (b_0^*, \dots, b_{d-1}^*),$$

with $\langle \phi | \in \mathcal{H}^*$.

Normalization of states:

$$\| |\psi\rangle \|^2 := \langle \psi | \psi \rangle = 1 \quad \implies \quad |\psi\rangle \sim e^{i\theta} |\psi\rangle$$

Orthonormal basis $\{|k\rangle\}_{k=0}^{d-1}$ on $\mathcal{H} = \mathbb{C}^d$:

$$\langle k | k' \rangle = \delta_{kk'} \implies |\psi\rangle = \sum_k c_k |k\rangle$$

Preliminaries

Schrödinger equation

Time evolution is governed by the **unitary operator** $U(t) : \mathcal{H} \rightarrow \mathcal{H}$:

$$\frac{d}{dt}|\psi(t)\rangle = -\frac{i}{\hbar}H|\psi(t)\rangle \iff |\psi(t)\rangle = U(t)|\psi(0)\rangle, \quad t \geq 0,$$

where $H = H^\dagger$ is the **Hamiltonian** of the system.

Preliminaries

Schrödinger equation

Time evolution is governed by the **unitary operator** $U(t) : \mathcal{H} \rightarrow \mathcal{H}$:

$$\frac{d}{dt}|\psi(t)\rangle = -\frac{i}{\hbar}H|\psi(t)\rangle \iff |\psi(t)\rangle = U(t)|\psi(0)\rangle, \quad t \geq 0,$$

where $H = H^\dagger$ is the **Hamiltonian** of the system.

Unitary property $U^\dagger(t)U(t) = U(t)U^\dagger(t) = \mathbb{1}$.

- Conservation of probability $\|U(t)\psi\|^2 = \|\psi\|^2$.
- Spectral theorem implies $U(t) = e^{-iHt/\hbar}$ (next slide).

Preliminaries

Spectral Theorem

Any **normal operator** $A : \mathcal{H} \rightarrow \mathcal{H}$ satisfying $A^\dagger A = A A^\dagger$ is diagonal with respect to some orthonormal basis $\{|a\rangle\}$:

$$A = \sum_{a \in \sigma(A)} a |a\rangle \langle a| \quad \implies \quad A |a\rangle = a |a\rangle,$$

where $\sigma(A)$ is the finite set of **eigenvalues** of A .

Preliminaries

Spectral Theorem

Any **normal operator** $A : \mathcal{H} \rightarrow \mathcal{H}$ satisfying $A^\dagger A = A A^\dagger$ is diagonal with respect to some orthonormal basis $\{|a\rangle\}$:

$$A = \sum_{a \in \sigma(A)} a |a\rangle \langle a| \quad \implies \quad A |a\rangle = a |a\rangle,$$

where $\sigma(A)$ is the finite set of **eigenvalues** of A .

- **Observables** are represented by self-adjoint operators $A = A^\dagger$, $a \in \mathbb{R}$.
- Hamiltonian $H = \sum_k E_k |k\rangle \langle k|$, with E_k the corresponding energies.

Preliminaries

Spectral Theorem

Any **normal operator** $A : \mathcal{H} \rightarrow \mathcal{H}$ satisfying $A^\dagger A = AA^\dagger$ is diagonal with respect to some orthonormal basis $\{|a\rangle\}$:

$$A = \sum_{a \in \sigma(A)} a |a\rangle \langle a| \quad \implies \quad A|a\rangle = a|a\rangle,$$

where $\sigma(A)$ is the finite set of **eigenvalues** of A .

- **Observables** are represented by self-adjoint operators $A = A^\dagger$, $a \in \mathbb{R}$.
- Hamiltonian $H = \sum_k E_k |k\rangle \langle k|$, with E_k the corresponding energies.

If $f : \mathbb{C} \rightarrow \mathbb{C}$ is an analytic function:

$$f(A) = \sum_a f(a) |a\rangle \langle a| \quad \implies \quad U(t) = \sum_k e^{-iE_k t/\hbar} |k\rangle \langle k|$$

Preliminaries

Projective measurements: For observable $A = \sum_a a|a\rangle\langle a|$, probability of measuring a is $p_a = |\langle a|\psi\rangle|^2$.

Expectation value $\langle A \rangle$:

$$\langle A \rangle = \sum_a a p_a = \sum_a a \langle a|\psi\rangle\langle\psi|a\rangle = \text{Tr}(A|\psi\rangle\langle\psi|)$$

Preliminaries

Projective measurements: For observable $A = \sum_a a|a\rangle\langle a|$, probability of measuring a is $p_a = |\langle a|\psi\rangle|^2$.

Expectation value $\langle A \rangle$:

$$\langle A \rangle = \sum_a a p_a = \sum_a a \langle a|\psi\rangle\langle\psi|a\rangle = \text{Tr}(A|\psi\rangle\langle\psi|)$$

Preliminaries

Projective measurements: For observable $A = \sum_a a|a\rangle\langle a|$, probability of measuring a is $p_a = |\langle a|\psi\rangle|^2$.

Expectation value $\langle A \rangle$:

$$\langle A \rangle = \sum_a a p_a = \sum_a a \langle a|\psi\rangle\langle\psi|a\rangle = \text{Tr}(A|\psi\rangle\langle\psi|)$$

Density matrix ρ

Pure state ensemble $\{p_{\psi_j}, |\psi_j\rangle\}$:

$$\langle A \rangle = \sum_a \sum_j a p_a p_{\psi_j} = \text{Tr}(A \sum_j p_{\psi_j} |\psi_j\rangle\langle\psi_j|),$$

where $\rho = \sum_j p_{\psi_j} |\psi_j\rangle\langle\psi_j|$ is the **density matrix**.

Properties:

- Unit trace: $\text{Tr}(\rho) = 1$.
- Hermitian: $\rho = \rho^\dagger$.
- Positive: $\langle\phi|\rho|\phi\rangle \geq 0, \quad \forall \phi \in \mathcal{H}$.

Dynamics

Liouville-von Neumann equation

Time evolution (Schrödinger equation):

$$\frac{d}{dt}\rho(t) = -i[H(t), \rho(t)] \iff \rho(t) = U(t)\rho U^\dagger(t), \quad t \geq 0,$$

where $U(t) = \mathcal{T}e^{-i \int_0^t ds H(s)}$ for time dependent $H(t)$.

Dynamics

Liouville-von Neumann equation

Time evolution (Schrödinger equation):

$$\frac{d}{dt}\rho(t) = -i[H(t), \rho(t)] \quad \Longleftrightarrow \quad \rho(t) = U(t)\rho U^\dagger(t), \quad t \geq 0,$$

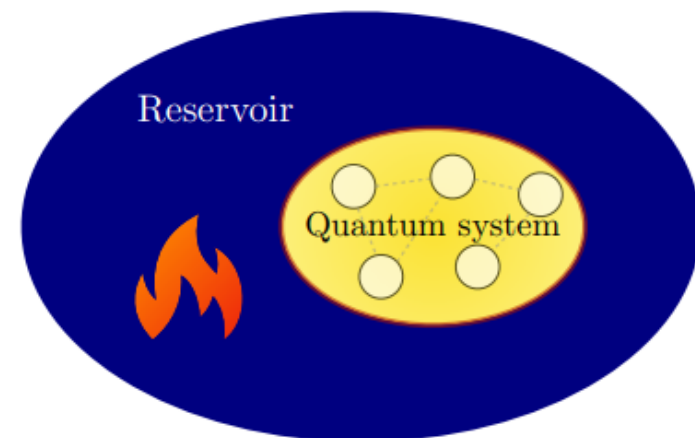
where $U(t) = \mathcal{T}e^{-i\int_0^t ds H(s)}$ for time dependent $H(t)$.

Derivation:

$$\begin{aligned} \frac{d}{dt}\rho(t) &= \left(\frac{d}{dt}U(t)\right)\rho U^\dagger(t) + U(t)\rho\left(\frac{d}{dt}U^\dagger(t)\right) \\ &= -\frac{i}{\hbar}H(t)\rho(t) + \frac{i}{\hbar}\rho(t)H(t) \\ &= -i[H(t), \rho(t)] \end{aligned}$$

Open setting

Open system S and reservoir R :



Open setting

Open system S and reservoir R :

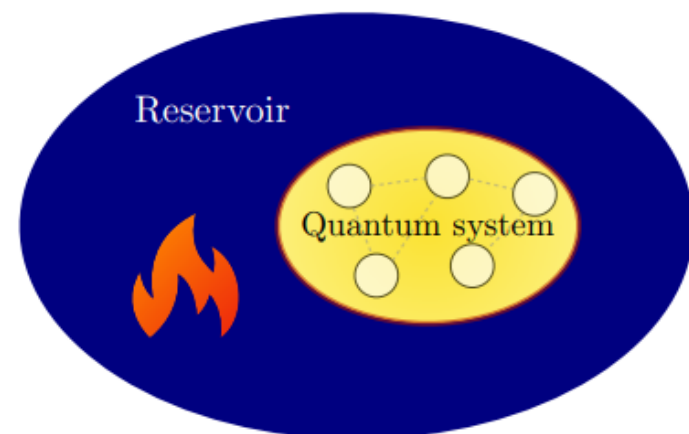
$$\mathcal{H} = \mathcal{H}_S \otimes \mathcal{H}_R$$

For orthonormal bases $\{|k\rangle_S\}$ and $\{|\mu\rangle_R\}$:

$$|\psi\rangle = \sum_{k,\mu} c_{k,\mu} |k\rangle_S \otimes |\mu\rangle_R, \quad \mathcal{H} = \text{span}\{|k\rangle_S \otimes |\mu\rangle_R\}$$

Density matrix:

$$\rho = \sum_a \sum_{j,k} \sum_{\mu,\nu} p_{\psi_a} c_{a,j\mu} c_{a,k\nu}^* |j\rangle \langle k|_S \otimes |\mu\rangle \langle \nu|_R$$



Open setting

Open system S and reservoir R :

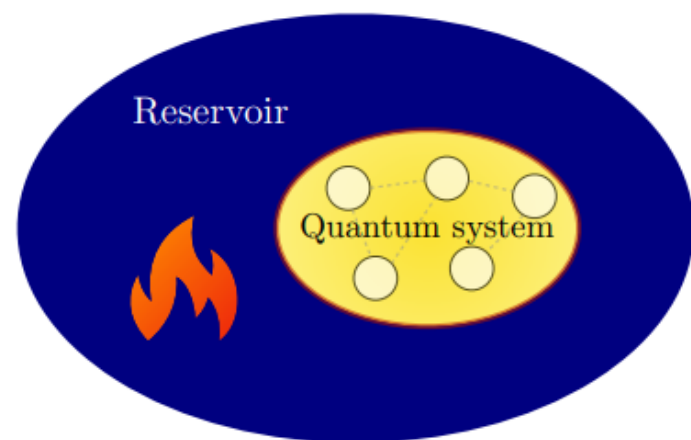
$$\mathcal{H} = \mathcal{H}_S \otimes \mathcal{H}_R$$

For orthonormal bases $\{|k\rangle_S\}$ and $\{|\mu\rangle_R\}$:

$$|\psi\rangle = \sum_{k,\mu} c_{k,\mu} |k\rangle_S \otimes |\mu\rangle_R, \quad \mathcal{H} = \text{span}\{|k\rangle_S \otimes |\mu\rangle_R\}$$

Density matrix:

$$\rho = \sum_a \sum_{j,k} \sum_{\mu,\nu} p_{\psi_a} c_{a,j\mu} c_{a,k\nu}^* |j\rangle \langle k|_S \otimes |\mu\rangle \langle \nu|_R$$



Reduced density matrix

Expectation value:

$$\langle A_S \otimes \mathbb{1}_R \rangle = \text{Tr}_S[(A_S \otimes \mathbb{1}_R) \text{Tr}_R(\rho)] = \text{Tr}(A_S \rho_S),$$

where $\rho_S = \text{Tr}_R(\rho) = \sum_{\mu} \langle \mu | \rho | \mu \rangle_R$ is the **reduced system density matrix**.

Open setting

Time evolution of open system S :

$$\rho_S(t) = \text{Tr}_R \left(U(t) \rho(0) U^\dagger(t) \right)$$

Open setting

Time evolution of open system S :

$$\rho_S(t) = \text{Tr}_R \left(U(t) \rho(0) U^\dagger(t) \right)$$

Separable $\rho = \rho_S \otimes \rho_R$:

$$\begin{aligned} \rho_S(t) &= \sum_{\mu} \langle \mu | U(t) \rho_S \otimes \sum_{\nu} \lambda_{\nu} |\nu\rangle \langle \nu| U^\dagger(t) | \mu \rangle \\ &= \sum_{\mu, \nu} \sqrt{\lambda_{\nu}} \langle \mu | U(t) | \nu \rangle \rho_S \sqrt{\lambda_{\nu}} \langle \nu | U^\dagger(t) | \mu \rangle = \sum_{\mu, \nu} K_{\mu\nu}(t) \rho_S(0) K_{\mu\nu}^\dagger(t) \end{aligned}$$

Open setting

Time evolution of open system S :

$$\rho_S(t) = \text{Tr}_R \left(U(t) \rho(0) U^\dagger(t) \right)$$

Separable $\rho = \rho_S \otimes \rho_R$:

$$\begin{aligned} \rho_S(t) &= \sum_{\mu} \langle \mu | U(t) \rho_S \otimes \sum_{\nu} \lambda_{\nu} |\nu\rangle \langle \nu | U^\dagger(t) | \mu \rangle \\ &= \sum_{\mu, \nu} \sqrt{\lambda_{\nu}} \langle \mu | U(t) | \nu \rangle \rho_S \sqrt{\lambda_{\nu}} \langle \nu | U^\dagger(t) | \mu \rangle = \sum_{\mu, \nu} K_{\mu\nu}(t) \rho_S(0) K_{\mu\nu}^\dagger(t) \end{aligned}$$

Kraus representation

The time evolution of an open quantum system under the **separability** condition can be described by a **quantum map** $\Phi_t : \mathcal{S}(\mathcal{H}_S) \rightarrow \mathcal{S}(\mathcal{H}_S)$ in **Kraus form**:

$$\rho_S(t) = \Phi_t \rho_S(0) = \sum_{\alpha=1}^{d^2} K_{\alpha}(t) \rho_S(0) K_{\alpha}^\dagger(t)$$

where $\sum_{\alpha} K_{\alpha}^\dagger K_{\alpha} = \mathbb{1}_S$.

K. Kraus, vol. **190** of Springer Lecture Notes in Physics. Springer-Verlag, Berlin, 1983

Equilibrium quantities

Von Neumann entropy of state ρ :

$$S(\rho) = -k_B \text{Tr}(\rho \ln \rho) = -k_B \sum_a \lambda_a \ln \lambda_a \geq 0$$

from the spectral decomposition $\rho = \sum_a \lambda_a |a\rangle \langle a|$.

Equilibrium quantities

Von Neumann entropy of state ρ :

$$S(\rho) = -k_B \text{Tr}(\rho \ln \rho) = -k_B \sum_a \lambda_a \ln \lambda_a \geq 0$$

from the spectral decomposition $\rho = \sum_a \lambda_a |a\rangle\langle a|$.

Gibbs state

Maximization of entropy:

$$S(\rho) = -k_B \text{Tr}(\rho \ln \rho) \leq -k_B \text{Tr}(\rho \ln \rho_G) = S_G \quad \implies \quad S(\rho) \leq S(\rho_G)$$

where

$$\rho_G = \frac{e^{-\beta H}}{Z}, \quad \beta = \frac{1}{k_B T}, \quad Z = \text{Tr} e^{-\beta H}$$

Gibbs state equivalent to **canonical ensemble**.

Relative entropy:

$$S(\rho \| \sigma) = \text{Tr}(\rho \ln \rho) - \text{Tr}(\rho \ln \sigma) \geq 0$$

Overview

In this lecture we have:

- Introduced quantum thermodynamics as an extension of macroscopic thermodynamics to the microscopic domain.
- Reviewed concepts of quantum probability and statistics.
- Formulated the time evolution of open quantum systems in terms of quantum maps.
- Reviewed equilibrium concepts.

Next lecture:

- Markovian master equations.
- 1st and 2nd laws of thermodynamics.

Thank you



quantum.sun.ac.za